

SOME INTEGRALS INVOLVING THREE BESSEL FUNCTIONS

by

R. K. Bhatnagar

Weapons Evaluation Group, Delhi-6

In this paper we shall evaluate the integrals of the following type:—

$$I = \int_0^1 X J_0(\lambda X) I_s(\mu X) I_s(\nu X) dX$$

$$J = \int_0^1 X J_0(\lambda X) J_s(\mu X) J_s(\nu X) dX$$

$$K = \int_0^1 X^\lambda J_\nu(bX) J_0(cX) J_0(dX) dX$$

which we have come across in the solution of some non-linear problems with central or axial symmetry.

Evidently these integrals provide Fourier-Bessel expansions for $I_s(\mu X) I_s(\nu X)$, $J_s(\mu X) J_s(\nu X)$ and $J_\nu(bX) J_0(cX)$ respectively if λ is one of the positive zeroes of the function $J_0(X)$. To the best knowledge of the author the values of these integrals are not available in the current literature.

1. Here we shall evaluate the following integral

$$I = \int_0^1 X J_0(\lambda X) I_s(\mu X) I_s(\nu X) dX \quad \dots \quad (1)$$

We shall use the following results:—

$$\int_0^1 X J_0(\lambda X) J_0(KX) dX = \frac{\lambda}{K^2 - \lambda^2} J_0'(\lambda) J_0(K) \quad \dots \quad (1.1)$$

and the addition theorem

$$J_0\left\{X\sqrt{\mu^2 + \nu^2 - 2\mu\nu \cos \phi}\right\} = \sum_{s=0}^{\infty} \epsilon_s J_s(\mu X) J_s(\nu X) \cos S\phi \quad \dots \quad (1.2)$$

where $\epsilon_0 = 1$, $\epsilon_s = 2$ for $S = 1, 2, \dots$

Taking $K = i \sqrt{\mu^2 + \nu^2 - 2\mu\nu \cos \phi}$ in (1.1) and integrating with respect to ϕ between 0 and π after multiplying by $\cos S\phi$ we get

$$\begin{aligned} & \int_0^1 X J_0(\lambda X) I_s(\mu X) I_s(\nu X) dX \\ &= \frac{(-1)^s \lambda J_1(\lambda)}{\pi \epsilon_s} \sum_{m=0}^{\infty} (-1)^m \epsilon_m I_m(\mu) I_m(\nu) \int_0^{\pi} \frac{\cos(m+s)\phi + \cos(m-s)\phi}{\lambda^2 + \mu^2 + \nu^2 - 2\mu\nu \cos \phi} d\phi \\ &= \frac{(-1)^s \lambda J_1(\lambda)}{\epsilon_s} \sum_{m=0}^{\infty} \frac{(-1)^m \epsilon_m I_m(\mu) I_m(\nu)}{(a^2 - b^2)^{1/2}} \left[\left\{ \frac{-\sqrt{a^2 - b^2} + a}{b} \right\}^{m+s} \right. \\ & \quad \left. + \left\{ \frac{-\sqrt{a^2 - b^2} + a}{b} \right\}^{m-s} \right] \end{aligned} \quad (1.3)$$

$$\text{where } a = \lambda^2 + \mu^2 + \nu^2, \quad b = 2\mu\nu \quad \dots \quad (1.4)$$

when $S=0$ we have the particular case

$$\begin{aligned} & \int_0^1 X J_0(\lambda X) I_0(\mu X) I_0(\nu X) dX \\ &= 2\lambda J_1(\lambda) \sum_{m=0}^{\infty} \frac{(-1)^m \epsilon_m I_m(\mu) I_m(\nu)}{[(\lambda^2 + \mu^2 + \nu^2)^2 - 4\mu^2\nu^2]^{1/2}} \\ & \quad \times \left\{ \frac{\lambda^2 + \mu^2 + \nu^2 - \sqrt{(\lambda^2 + \mu^2 + \nu^2)^2 - 4\mu^2\nu^2}}{2\mu\nu} \right\}^m \end{aligned} \quad (1.5)$$

2. We shall now evaluate the integral

$$J = \int_0^1 X J_0(\lambda X) J_s(\mu X) J_s(\nu X) dX \quad \dots \quad (2.1)$$

We know that

$$I_\nu(z) = e^{-\frac{1}{2}\pi i \nu} J_\nu\left(z e^{\frac{1}{2}\pi i}\right) \quad \dots \quad (2.2)$$

where $-\pi < \arg Z \leq \frac{1}{2}\pi$

$$\text{or } I_\nu\left(z e^{-\frac{1}{2}\pi i}\right) = e^{-\frac{1}{2}\pi i \nu} J_\nu(z)$$

Hence by replacing μ by $\mu e^{-\frac{1}{2}\pi i}$ and ν by $\nu e^{-\frac{1}{2}\pi i}$ we get

$$\int_0^1 X J_0(\lambda X) J_s(\mu X) J_s(\nu X) dX = \frac{\lambda J_0'(\lambda)}{\epsilon_s} \sum_{m=0}^{\infty} \epsilon_m J_m(\mu) J_m(\nu) B \quad \dots \quad (2.3)$$

where

$$B = \frac{1}{\pi} \int_0^{\pi} \frac{\cos(m+s)\phi + \cos(m \sim s)\phi}{\mu^2 + \nu^2 - \lambda^2 - 2\mu\nu \cos \phi} d\phi \quad \dots \quad (2.4)$$

The value of B depends on the relative values of λ , μ and ν .

when $\mu \sim \nu > \lambda$,

$$B = \frac{1}{(a^2 - b^2)^{\frac{1}{2}}} \left[\left\{ \frac{a - \sqrt{a^2 - b^2}}{b} \right\}^{m+s} + \left\{ \frac{a - \sqrt{a^2 - b^2}}{b} \right\}^{m \sim s} \right] \quad (2.5)$$

where

$$a = \mu^2 + \nu^2 - \lambda^2 > 0, \quad b = 2\mu\nu$$

when $\mu + \nu < \lambda$

$$B = - \frac{1}{(\alpha^2 - b^2)^{\frac{1}{2}}} \left[\left\{ \frac{\sqrt{\alpha^2 - b^2} - a}{b} \right\}^{m+s} + \left\{ \frac{\sqrt{\alpha^2 - b^2} - a}{b} \right\}^{m \sim s} \right] \quad (2.6)$$

$$\text{where } \alpha = \lambda^2 - \mu^2 - \nu^2 > 0$$

when $\mu \sim \nu < \lambda$ $\sqrt{\mu^2 + \nu^2}$ and $\sqrt{\mu^2 + \nu^2} < \lambda < \mu + \nu$ the integral becomes unbounded in the range of integration. In these two cases we shall take the principal values which exist.

when $\mu \sim \nu < \lambda < \sqrt{\mu^2 + \nu^2}$

$$\begin{aligned} B &= \frac{1}{2\pi\mu\nu} \int_0^{\pi} \frac{\cos(m+s)\phi + \cos(m \sim s)\phi}{\cos \psi - \cos \phi} d\phi \\ &= - \frac{1}{2\mu\nu} \left[\frac{\sin(m+s)\psi + \sin(m \sim s)\psi}{\sin \psi} \right] \quad \dots \quad (2.7) \end{aligned}$$

where

$$0 < \cos \psi = \frac{\mu^2 + \nu^2 - \lambda^2}{2\mu\nu} < 1 \quad \dots \quad (2.8)$$

so that $0 < \psi < \frac{\pi}{2}$

when $\sqrt{\mu^2 + \nu^2} < \lambda < \mu + \nu$

$$\begin{aligned} B &= - \frac{1}{2\pi\mu\nu} \int_0^{\pi} \frac{\cos(m+s)\phi + \cos(m \sim s)\phi}{\cos \phi - \cos \psi} d\phi \\ &= - \frac{1}{2\mu\nu} \left[\frac{\sin(m+s)\psi + \sin(m \sim s)\psi}{\sin \psi} \right] \quad \dots \quad (2.9) \end{aligned}$$

where $-1 < \cos \psi = \frac{\mu^2 + \nu^2 - \lambda^2}{2\mu\nu} < 0$

so that $\frac{\pi}{2} < \psi < \pi$

In particular when $S = 0$

$$\int_0^1 X J_0(\lambda X) J_0(\mu X) J_0(\nu X) dX = \lambda J_0'(\lambda) \sum_{m=0}^{\infty} \epsilon_m J_m(\mu) J_m(\nu) B \quad (2 \cdot 10)$$

Where

$$\begin{aligned} B &= \frac{2}{(a^2 - b^2)^{\frac{1}{2}}} \left\{ \frac{a - \sqrt{a^2 - b^2}}{b} \right\}^m, \mu \sim \nu < \lambda \\ &= -\frac{2}{(a^2 - b^2)^{\frac{1}{2}}} \left\{ \frac{\sqrt{a^2 - b^2} - a}{b} \right\}^m, \lambda < \mu + \nu \\ &= -\frac{1}{\mu\nu} \frac{\sin m\psi}{\sin \psi}, 0 < \psi < \frac{\pi}{2}, \mu \sim \nu < \lambda < \sqrt{\mu^2 + \nu^2} \\ &= -\frac{1}{\mu\nu} \frac{\sin m\psi}{\sin \psi}, \frac{\pi}{2} < \psi < \pi, \sqrt{\mu^2 + \nu^2} < \lambda < \mu + \nu \end{aligned} \quad (2 \cdot 11)$$

The integrals (1.5) and (2.11) provide Fourier-Bessel expansions for $I_0(\mu X)$, $I_0(\nu X)$ and $J_0(\mu X) J_0(\nu X)$ respectively in terms of $J_0(\lambda_n X)$ where λ_n 's are positive zeroes of $J_0(X)$.

3. In this section we shall evaluate the integral

$$K = \int_0^1 X^\lambda J_\nu(bX) J_0(cX) J_0(dX) dX \quad \dots \quad (3 \cdot 1)$$

From the multiplication formula (1)

$$\begin{aligned} J_\mu(aX) J_\nu(bX) &= \frac{(\frac{1}{2}az)^\mu (\frac{1}{2}bz)^\nu}{|\nu + 1|} \\ &\cdot \sum_{m=0}^{\infty} (-1)^m \frac{(\frac{1}{2}az)^{2m} {}_2F_1(-m, -\mu - m; \nu + 1; b^2/a^2)}{|m| |(\mu + m + 1)|} \end{aligned} \quad (3 \cdot 2)$$

We have

$$\begin{aligned} \int_0^1 X^\lambda J_\mu(aX) J_\nu(bX) dX &= \frac{(\frac{1}{2}a)^\mu (\frac{1}{2}b)^\nu}{|\nu + 1|} \sum_{m=0}^{\infty} (-1)^m \\ &\frac{(\frac{1}{2}a)^{2m} {}_2F_1(-m, -\mu - m, \nu + 1; b^2/a^2)}{|m| |(\mu + m + 1)| (\lambda + \mu + \nu + 2m + 1)} \end{aligned} \quad (3 \cdot 3)$$

and taking $\mu = \frac{1}{2}, \lambda = \frac{3}{2}$ we get

$$\int_0^1 X \sin aX J_\nu(bX) dX = \frac{(\frac{1}{2}a)^{\frac{1}{2}} (\frac{1}{2}b)^\nu}{|\nu+1|} \sum_{m=0}^{\infty} (-1)^m \frac{(\frac{1}{2}a)^{2m} {}_2F_1(-m, -\frac{1}{2}-m; \nu+1; b^2/a^2)}{|m| (m+\frac{3}{2}) (\nu+2m+3)} \quad (3.4)$$

and taking $\mu = -\frac{1}{2}, \lambda = \frac{3}{2}$, we get

$$\int_0^1 X \cos aX J_\nu(bX) dX = \frac{(\frac{1}{2}a)^{-\frac{1}{2}} (\frac{1}{2}b)^\nu}{|\nu+1|} \sum_{m=0}^{\infty} (-1)^m \frac{(\frac{1}{2}a)^{2m} {}_2F_1(-m, \frac{1}{2}-m; \nu+1, b^2/a^2)}{|m| (m+\frac{1}{2}) (\nu+2m+2)} \quad (3.5)$$

Let $\mu = 0, a = \sqrt{c^2 + d^2 - 2cd \cos \phi}$ in (3.3) (3.6)

Integrating between 0 and π with respect to ϕ we get

$$\int_0^1 X \lambda J_\nu(bX) J_0(cX) J_0(dX) dX = 2 \frac{(\frac{1}{2}b)^\nu}{|\nu+1|} \sum_{m=0}^{\infty} (-1)^m \frac{1}{2^{2m} (|m|)^2 |\nu+\lambda+2m+1|} \quad (3.7)$$

where

$$\begin{aligned} I &= \frac{1}{\pi} \int_0^\pi (c^2 + d^2 - 2cd \cos \phi)^m {}_2F_1\left(-m, -m; \nu+1, \frac{b^2}{c^2 + d^2 - 2cd \cos \phi}\right) d\phi \\ &= \frac{1}{\pi} \sum_{r=0}^m \frac{(-m)_r (-m)_r}{(\nu+1)_r r!} b^{2r} \int_0^\pi (c^2 + d^2 - 2cd \cos \phi)^{m-r} d\phi \\ &= \sum_{r=0}^m \frac{(-m)_r (-m)_r}{(\nu+1)_r r!} b^{2r} (c^2 \sim d^2)^{m-r} P_{m-r}\left(\frac{c^2 + d^2}{c^2 \sim d^2}\right) \quad (3.8) \end{aligned}$$

where $P_n(z)$ are usual Legendre's Polynomials.

Acknowledgement

The author is grateful to Dr. R. S. Varma, Director, Defence Science Laboratory for his keen interest and Prof. S.P. Chakravarti, Director, Weapons Evaluation Group for his encouragement in the preparation of this paper.

References

1. Watson, G. N.,—A Treatise on the Theory of Bessel Functions, 148, 1922.
2. Whittaker, E. T. and Watson, G. N.,—A Course on Modern Analysis, 312, 1952.