

FOURIER SERIES FOR FOX'S H -FUNCTION OF TWO VARIABLES

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An attempt has been made to derive a Fourier series expansion for the H -function of two variables recently defined by Verma. This series is analogous to that of other special functions such as the MacRobert's E -function, Meijer's G -function and Fox's H -function of single variable as given by MacRobert, Kesarwani, Parihar, Parashar, Kapoor & Gupta. In the end an integral has been evaluated by making use of this result.

MacRobert¹, Kesarwani², Parihar³ established the Fourier series for the E - and G -function and in the recent paper Parashar⁴, Kapoor & Gupta⁵ has proved Fourier series for Fox's H -function of single variable. However the Fourier series expansion for H -function of two variables has not been derived so far.

The following Fourier series expansion is proposed to be established :

$$\sum_{r=0}^{\infty} \frac{(k+r)!}{k! r!} H_{p, (t+3:t'), s, (q+3:q')}^{n, \nu_1+1, \nu_2, m_1+2, m_2} \left[x \begin{cases} (a_p, e_p) \\ (r, h), (\gamma_t, c_t), (0, h), (-k-r-1, h); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s, d_s) \end{cases} \right] \sin(k+2r+1)\theta \\ = \frac{\sqrt{-\pi}}{2} \sum_{u=0}^k \frac{\sin \theta (\cos \theta - 1)}{u! (k-u)!} H_{p, (t+2:t'), s, (q+2:q')}^{n, \nu_1, \nu_2, m_1+2, m_2} \left[\begin{array}{l} \frac{x}{\sin^{2h} \theta} \begin{cases} (a_p, e_p) \\ (\gamma_t, c_t), (0, h), (-\frac{1}{2}u, h); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s, d_s) \end{cases} \\ y \begin{cases} \left(1 + \frac{k}{2}, h\right), \left(\frac{3}{2} + \frac{k}{2}, h\right), (\beta_q, b_q), (1, h); (\beta'_{q'}, b'_{q'}) \\ \left(1 + \frac{k+u}{2}, h\right), \left(\frac{3}{2} + \frac{k+u}{2}, h\right), (\beta_q, b_q); \\ (\beta'_{q'}, b'_{q'}) \end{cases} \end{array} \right] \quad (1)$$

where $0 \leq \theta \leq \pi$, $T = \sum_1^n e_j + \sum_1^{\nu_1} c_j + \sum_1^{m_1} b_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{v_1+1}^t c_j - \sum_{m_1+1}^q b_j > 0$, $|\arg x| < \frac{1}{2} T \pi$,

and

$$T' = \sum_1^n e_j + \sum_1^{\nu_2} c'_j + \sum_1^{m_2} b'_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{v_1+1}^{t'} c'_j - \sum_{m_2+1}^{q'} b'_j > 0, |\arg y| < \frac{1}{2} T' \pi.$$

Fox's H -function of two variables recently introduced by Verma⁶ which is an extension of G -function of two variables defined by Agarwal⁷. This H -function of two variables does not only includes Fox's H -function and the Meijer's G -function of single variables as particular cases but also most of special functions of two variables, e.g., Appell's functions, the Whittaker function of two variables etc.

Thus Fox's H -function of two variables due to Verma⁶ will be defined as follows:

$$H_{p, (t:t'), s, (q:q')}^{n, \nu_1, \nu_2, m_1, m_2} \left[\begin{array}{l} x \begin{cases} (a_p, e_p) \\ (\gamma_t, c_t); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s, d_s) \end{cases} \\ y \begin{cases} (\beta_q, b_q); (\beta'_{q'}, b'_{q'}) \end{cases} \end{array} \right] = \frac{1}{(2\pi i)^2} \int_{-i\infty}^{i\infty} \int_{-i\infty}^{i\infty} \phi(\xi + \eta) \psi(\xi, \eta) x^\xi y^\eta d\xi d\eta, \quad (2)$$

where

$$\phi(\xi + \eta) = \frac{\prod_{j=1}^n \Gamma(1 - a_j + e_j \xi + e_j \eta)}{\prod_{j=n+1}^p \Gamma(a_j - e_j \xi - e_j \eta) \prod_{j=1}^s \Gamma(\delta_j + d_j \xi + d_j \eta)},$$

$$\psi(\xi, \eta) = \frac{\prod_{j=1}^{m_1} \Gamma(\beta_j - b_j \xi) \prod_{j=1}^{\nu_1} \Gamma(\gamma_j + c_j \xi) \prod_{j=1}^{m_2} \Gamma(\beta'_j - b'_j \eta) \prod_{j=1}^{\nu_2} \Gamma(\gamma'_j + c'_j \eta)}{\prod_{j=m_1+1}^q \Gamma(1 - \beta_j + b_j \xi) \prod_{j=\nu_1+1}^t \Gamma(1 - \gamma_j - c_j \xi) \prod_{j=m_2+1}^{q'} \Gamma(1 - \beta'_j + b'_j \eta) \prod_{j=\nu_2+1}^{t'} \Gamma(1 - \gamma'_j - c'_j \eta)},$$

and

$$0 \leq m_1 \leq q, 0 \leq m_2 \leq q', 0 \leq \nu_1 \leq t, 0 \leq \nu_2 \leq t', 0 \leq n \leq p.$$

The sequence of parameters $(\beta_{m_1}, b_{m_1}), (\beta'_{m_2}, b'_{m_2}), (\gamma_{\nu_1}, c_{\nu_1}), (\gamma'_{\nu_2}, c'_{\nu_2})$ and (a_n, e_n) are such that none of the poles of integrand coincides. The paths of integration are indented, if necessary, in such a manner that all the poles of $\Gamma(\beta_j - b_j \xi), j = 1, 2, \dots, m_1$ and $\Gamma(\beta'_k - b'_k \eta), k = 1, 2, \dots, m_2$ lie to the right and those of $\Gamma(\gamma_j + c_j \xi), j = 1, 2, \dots, \nu_1$, $\Gamma(\gamma'_k + c'_k \eta), k = 1, 2, \dots, \nu_2$ and $\Gamma(1 - a_j + e_j \xi + e_j \eta), j = 1, 2, \dots, n$, lie to the left of imaginary axis.

The integral (2) converges if

$$T = \sum_1^n e_j + \sum_1^{\nu_1} c_j + \sum_1^{m_1} b_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{\nu_1+1}^t c_j - \sum_{m_1+1}^q b_j > 0, |\arg x| < \frac{1}{2} T \pi,$$

and

$$T^* = \sum_1^n e_j + \sum_1^{\nu_2} c_j + \sum_1^{m_2} b'_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{\nu_2+1}^{t'} c'_j - \sum_{m_2+1}^{q'} b'_j > 0, |\arg y| < \frac{1}{2} T^* \pi$$

We shall give below some results left and use them later on. Askey⁸ give with $\lambda = 1-s$

$$(\sin \theta)^{1-2s} P_n^{1-s}(\cos \theta) = \sum_{r=0}^{\infty} \frac{2^{2s}(n+r)!}{\Gamma(1-s)\Gamma(s)r!n!} \frac{\Gamma(n+2-2s)\Gamma(r+s)}{\Gamma(n+r+2-s)} \sin(n+2r+1)\theta, \quad (3)$$

where $s < 1$ and $0 \leq \theta \leq \pi$ and $P_n^{(\lambda)}(\cos \theta)$ is given by

$$(1 - 2r \cos \theta + r^2)^{-\lambda} = \sum_{n=0}^{\infty} P_n^{(\lambda)}(\cos \theta) \cdot r^n, \quad (4)$$

also Rainville⁹

$$P_n^{(\lambda)}(z) = \sum_{m=0}^n \frac{(2\lambda)_{m+n} \left(\frac{z-1}{2}\right)^m}{m!(n-m)! (\lambda+\frac{1}{2})_m}. \quad (5)$$

The Legendre duplication formula

$$\sqrt(\pi) \Gamma(2z) = 2^{2z-1} \Gamma(z) \Gamma(z + \frac{1}{2}). \quad (6)$$

Verma⁶ gives

$$H_{p,(t:t'),s,(q:q')}^{n,\nu_1,\nu_2,m_1,m_2} \begin{Bmatrix} x \\ y \end{Bmatrix} \begin{pmatrix} (a_p, e_p) \\ (\gamma_t, c_t); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s, d_s) \\ (\beta_q, b_q); (\beta'_{q'}, b'_{q'}) \end{pmatrix} = x^{-r} H_{p,(t:t'),s,(q:q')}^{n,\nu_1,\nu_2,m_1,m_2} \begin{Bmatrix} x \\ y \end{Bmatrix} \begin{pmatrix} (a_p + e_p r, e_p) \\ (\gamma_t - c_t r, c_t); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s - d_s r, d_s) \\ (\beta_q + b_q r, b_q); (\beta'_{q'}, b'_{q'}) \end{pmatrix} \quad (7)$$

and

$$H_{p,(t:t'),s,(q:q')}^{n,\nu_1,\nu_2,m_1,m_2} \begin{Bmatrix} x \\ y \end{Bmatrix} \begin{pmatrix} (a_p, 1) \\ (\gamma_t, 1); (\gamma'_{t'}, 1) \\ (\delta_s, 1) \\ (\beta_q, 1); (\beta'_{q'}, 1) \end{pmatrix} = G_{p,(t:t'),s,(q:q')}^{n,\nu_1,\nu_2,m_1,m_2} \begin{Bmatrix} x \\ y \end{Bmatrix} \begin{pmatrix} (a_p) \\ (\gamma_t); (\gamma'_{t'}) \\ (\delta_s) \\ (\beta_q); (\beta'_{q'}) \end{pmatrix}. \quad (8)$$

Proof of equation (I) : On expressing the H-function of two variables as Mellin-Barnes type of double integral in L.H.S. of (1) and changing the order of summation and integration as permissible by absolute convergence for stated conditions in (1), the series becomes

$$\frac{1}{(2\pi i)^2} \int_{-\infty}^{i\infty} \int_{-\infty}^{i\infty} \phi(\xi + \eta) \psi(\xi, \eta) \left[\sum_{r=0}^{\infty} \frac{(k+r)!}{k! r! \Gamma(h\xi)} \frac{\Gamma\left(1 + \frac{k}{2} - h\xi\right)}{\Gamma(1-h\xi)} \frac{\Gamma\left(\frac{3}{2} + \frac{k}{2} - h\xi\right)}{\Gamma(k+r+2-h\xi)} \right. \\ \left. \cdot \sin(k+2r+1)\theta \right] x^\xi y^\eta d\xi d\eta.$$

Using the result (3) and (6), we have

$$\frac{1}{(2\pi i)^2} \int_{-\infty}^{i\infty} \int_{-\infty}^{i\infty} \phi(\xi + \eta) \psi(\xi, \eta) \left[\frac{\sqrt{(\pi)}}{2^{\frac{k+1}{2}}} (\sin \theta)^{(1-2h)} P_k^{(1-h\xi)} (\cos \theta) \right] x^\xi y^\eta d\xi d\eta.$$

Now substituting the value of $P_k^{(1-h\xi)} (\cos \theta)$ from (5) and then using the result (6), we get

$$\sum_{u=0}^k \frac{\sqrt{(\pi)}}{2} \frac{\sin \theta (\cos \theta - 1)^u}{u! (k-u)!} \left[\frac{1}{(2\pi i)^2} \int_{-\infty}^{i\infty} \int_{-\infty}^{i\infty} \phi(\xi + \eta) \psi(\xi, \eta) \frac{\Gamma\left(1 + \frac{k+u}{2} - h\xi\right)}{\Gamma(1-h\xi)} \frac{\Gamma\left(\frac{3}{2} + \frac{k+u}{2} - h\xi\right)}{\Gamma\left(\frac{3}{2} + u - h\xi\right)} \right. \\ \left. \cdot \left(\frac{x}{\sin^{2h} \theta} \right)^\xi y^\eta d\xi d\eta \right],$$

by definition of H-function of two variables (2), we get R.H.S. of (1), which completes the proof.

We shall derive here other Fourier series by applying the property of the H-function or by specializing the parameters.

(i). The following Fourier series for H-function is arrived by using the result (7) :

$$\sum_{r=0}^{\infty} \frac{(k+r)!}{k! r!} H_{p, (t+3:t'), s, (q+3:q')}^{n, v_1+1, v_2, m_1+2, m_2} \left[\begin{array}{l} x \left| \begin{array}{l} (a_p + e_p r, e_p) \\ (r-hr, h), (y_t - c_t r, c_t), (-hr, h), (-k-r-hr-1, h); \\ (y'_t, c'_t) \end{array} \right. \\ y \left| \begin{array}{l} (\delta_s - d_s r, d_s) \\ \left(1 + \frac{k}{2} + hr, h\right), \left(\frac{3}{2} + \frac{k}{2} + hr, h\right), (\beta_q + \alpha_q r, b_q), (1 + hr, h); \\ (\beta'_q, b'_q) \end{array} \right. \end{array} \right].$$

$$x^{-r} \sin(k+2r+1)\theta = \frac{\sqrt{(\pi)}}{2} \sum_{u=0}^k \frac{\sin \theta (\cos \theta - 1)^u}{u! (k-u)!} H_{p, (t+2:t'), s, (q+2:q')}^{n, v_1, v_2, m_1+2, m_2}$$

$$\left[\begin{array}{l} x \left| \begin{array}{l} (a_p, e_p) \\ (y_t, c_t), (0, h), (-\frac{1}{2} - u, h); (y'_t, c'_t) \\ (\delta_s, d_s) \end{array} \right. \\ y \left| \begin{array}{l} \left(1 + \frac{k+u}{2}, h\right), \left(\frac{3}{2} + \frac{k+u}{2}, h\right), (\beta_q, b_q); (\beta'_q, b'_q) \end{array} \right. \end{array} \right] \quad (9)$$

(ii) On putting $h = 1 = (e_p) = (c_t) = (c'_{t'}) = (d_s) = (b_q) = (b'_{q'})$ in (1), and using (8), we get the following result for Meijer G-function of two variables :

$$\sum_{r=0}^{\infty} \frac{(k+r)!}{k! r!} G_p^{n, v_1+1, v_2, m_1+2, m_2, (t+3:t'), s, (q+3:q')} \left[\begin{array}{l} (a_p) \\ x \left| \begin{array}{l} r, (\gamma_t), 0, -k-r-1; (\gamma'_{t'}) \\ (\delta_s) \end{array} \right. \\ y \left| \begin{array}{l} k, 3 \\ 1 + \frac{k}{2}, \frac{3}{2} + \frac{k}{2}, (\beta_q), 1; (\beta'_{q'}) \end{array} \right. \end{array} \right] \sin(k+2r+1) \theta \\ = \frac{\sqrt{(\pi)}}{2} \sum_{u=0}^k \frac{\sin \theta (\cos \theta - 1)^u}{u! (k-u)!} G_p^{n, v_1, v_2, m_1+2, m_2, (t+2:t'), s, (q+2:q')} \left[\begin{array}{l} (a_p) \\ \frac{x}{\sin^2 \theta} \left| \begin{array}{l} (\gamma_t), 0, -\frac{1}{2}-u; (\gamma'_{t'}) \\ (\delta_s) \end{array} \right. \\ y \left| \begin{array}{l} k+u, 3 \\ 1 + \frac{k+u}{2}, \frac{3}{2} + \frac{k+u}{2}, (\beta_q), (\beta'_{q'}) \end{array} \right. \end{array} \right]. \quad (10)$$

From (1), we can easily deduce the following integral :

$$\int_0^{\pi} \sum_{u=0}^k H_p^{n, v_1, v_2, m_1+2, m_2, (t+2:t'), s, (q+2:q')} \left[\begin{array}{l} (a_p, e_p) \\ \frac{x}{\sin^2 h \theta} \left| \begin{array}{l} (\gamma_t, c_t), (0, h), (-\frac{1}{2}-u, h); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s d_s) \end{array} \right. \\ y \left| \begin{array}{l} k+u, h \\ 1 + \frac{k+u}{2}, h, \left(\frac{3}{2} + \frac{k+u}{2}, h \right), (\beta_q, b_q); (\beta'_{q'}, b'_{q'}) \end{array} \right. \end{array} \right] \\ \cdot \frac{\sin \theta \cdot \sin(k+2r+1) \theta \cdot (\cos \theta - 1)^u}{u! (k-u)!} d\theta = \sum_{r=0}^{\infty} \frac{\sqrt{(\pi)} (k+r)!}{k! r!} H_p^{n, v_1+1, v_2, m_1+2, m_2, (t+3:t'), s, (q+3:q')} \\ \left[\begin{array}{l} (a_p, e_p) \\ x \left| \begin{array}{l} (r, h), (\gamma_t, c_t), (0, h), (-k-r-1, h); (\gamma'_{t'}, c'_{t'}) \\ (\delta_s, d_s) \end{array} \right. \\ y \left| \begin{array}{l} k, h \\ 1 + \frac{k}{2}, h, \left(\frac{3}{2} + \frac{k}{2}, h \right), (\beta_q, b_q), (1, h); (\beta'_{q'}, b'_{q'}) \end{array} \right. \end{array} \right], \quad (11)$$

where $0 \leq \theta \leq \pi$, $T = \sum_1^n e_j + \sum_1^{v_1} c_j + \sum_1^{m_1} b_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{v_1+1}^t c_j - \sum_{m_1+1}^q b_j > 0$, $|\arg x| < \frac{1}{2} T \pi$,

and

$$T' = \sum_1^n e_j + \sum_1^{v_2} c'_j + \sum_1^{m_2} b'_j - \sum_{n+1}^p e_j - \sum_1^s d_j - \sum_{v_2+1}^{t'} c'_j - \sum_{m_2+1}^{q'} b'_j > 0, |\arg y| < \frac{1}{2} T' \pi.$$

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