

## ROCKET FLAMES

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Among the important parameters which characterise the rocket flames are the (1) velocity, (2) pressure, (3) temperature of the exhaust gases and (4) the nature of chemical reactions in the flame of such gases. For the determination of these quantities ordinary methods fail because the flow of exhaust gases is supersonic in character. An introduction of a probe or any foreign body will create such strong disturbances in the supersonic flow that the readings of observing instruments will have no value. Spectroscopic methods are therefore eminently suitable because observations can be taken on the flame under running conditions.

Determination of velocity by spectroscopic method is based upon the utilization of the well known Doppler effect. For a body moving with a velocity  $v$ , towards or away from an observer, making an angle  $\theta$  with the line of sight, the displacement of wave length is—

$$\frac{\delta\lambda}{\lambda} = \pm \frac{v}{c} \cos \theta$$

where  $c$  is the velocity of light. If the observations are taken in two different directions  $\theta_1, \theta_2$  the difference in the two displacements will be

$$\frac{\delta\lambda}{\lambda} = \frac{v}{c} (\cos \theta_1 - \cos \theta_2)$$

For any velocity attainable on the earth  $v/c$  is so small that it is difficult to detect the shift. The stars, however, show appreciable shifts since they move with velocities ranging between 10 and 30 Km/Sec, with some as high as 300 Km/Sec.

In rockets, the velocities of exhaust gases are about 2 Km/Sec, giving a spectral shift of less than 0.1 Angstrom unit. For detection of such small shifts, instruments of highest resolving power have to be used and the one recently used is the Febry-Perot Interferometer, whose principle is as follows :

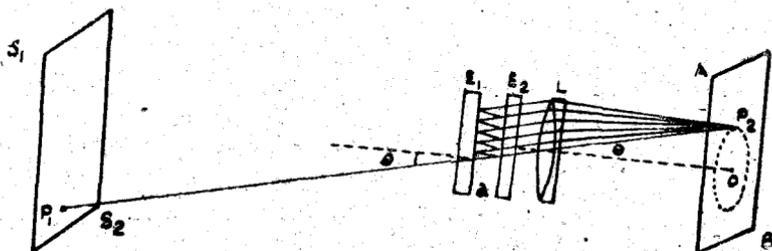


FIG. 1

\* Paper read on April 24, 1952 at the Defence Science Conference held in Delhi.

$E_1, E_2$ , are two glass plates thinly silvered on the inside. Light from a broad source  $S_1, S_2$  suffers multiple reflections and circular fringes are produced on the screen  $AB$ . A ray from a point  $P_1$  on the source is incident on the plate at an angle  $\theta$ ; the ray suffers several reflections between the two plates and the parallel reflections are brought to the point  $P_2$ . The condition for reinforcement is  $2d \cos \theta = m\lambda$ . The wave-length interval between successive fringes for  $\lambda = 5893 \times 10^{-8}$  cms, with the interferometer set at  $d = 0.2034$  cm is  $0.85 \text{ A}^\circ$ . If velocity of the exhaust gases is of the order of 2 Km/Sec.,  $\delta\lambda$  is of the order  $0.06 \text{ A}^\circ$ ; thus the Doppler shift is less than one tenth of the fringe setting. Before these principles are applied to the determination of the velocity, it is necessary to consider briefly the characteristics of the different zones in the exhaust flames. These are shown in fig. 2.

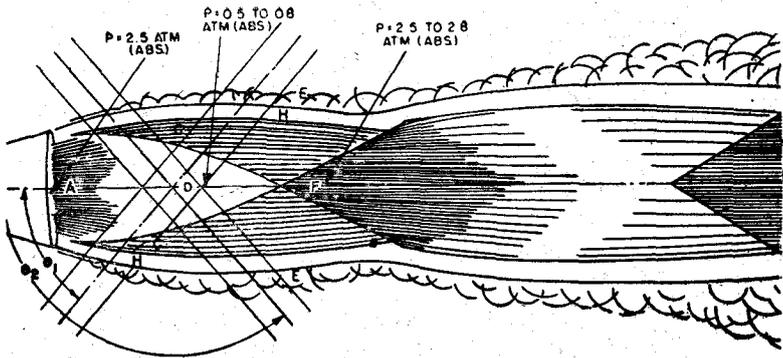


FIG. 2

If the line of sight is  $\theta_1$ , the flame zones across it are, the thin subsonic boundary layer  $E$ ; the relatively cool faintly luminous supersonic zone  $H$ , the hot supersonic zone  $C$ , the relatively cool overexpanded faintly luminous supersonic zone  $D$ . These are followed by repeated sections of zones  $C, H$  &  $E$ . Spectroscopically the two sections of zone  $C$  contribute the most light, zone  $D$  considerably less and zones  $H$  and  $E$  very little. The distribution of mass flow along the above zone is ——— one fourth in zone  $H$ , less than half in zone  $C$  and about one third in zone  $D$ . The latter zone records highest velocity owing to overexpansion, while owing to cooling and friction by the walls, the velocity in zone  $H$  is lower than the average.

The experimental arrangements for the determination of velocities of the exhaust gases are shown in the fig. 3.

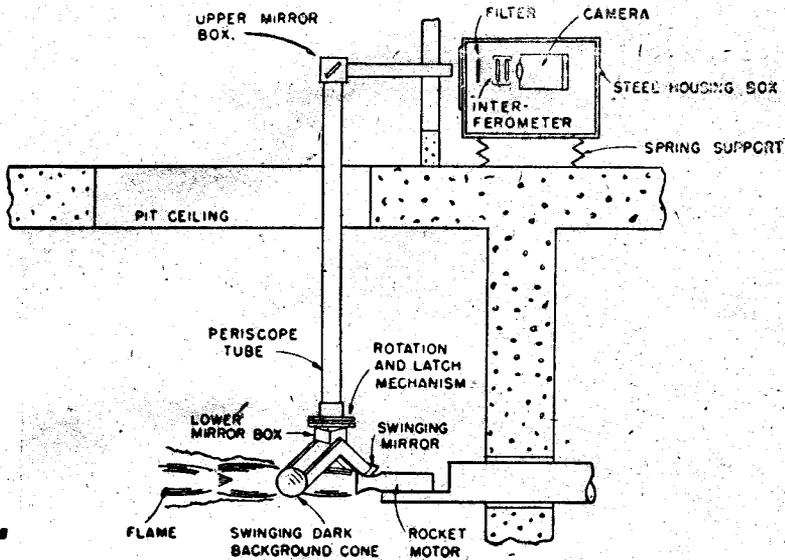


Diagram of the optical setup at the rocket test stand.

FIG. 3

Since there are strong vibrations near the rocket, the Fabry-Perot interferometer and the Camera are placed on a shock proof stand away from the rocket. Real image of the flame is focussed by the lens system of a periscope. For the velocity measurements, the periscope mirrors are so arranged that the same portion of the flame are viewed from two different angles  $\theta_1$  and  $\theta_2$  from the upstream flame axis. The Doppler shift of wave-length between these two angles of view is given by the relation,

$$\frac{\delta\lambda}{\lambda} = \frac{v}{c} (\cos \theta_1 - \cos \theta_2)$$

where  $v$  and  $c$  are the velocities of the gas and of light respectively. The interference patterns, corresponding to the two angles of view  $\theta_1$  and  $\theta_2$  are shown in fig. 4.

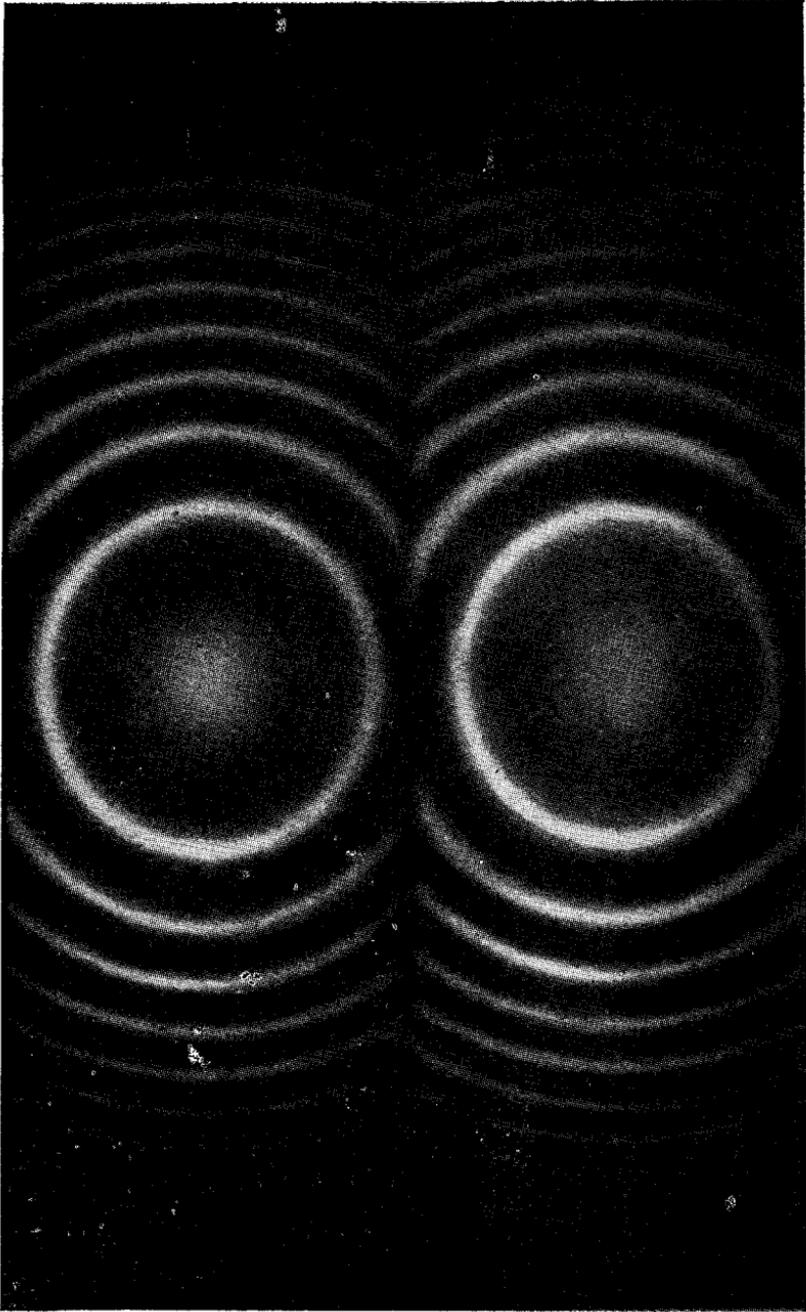


FIG. 4



The figures for the velocity, pressure and temperature for two different regions, D and F (See fig. 2) are

		D	F
Temperature	.. .. .	2200°K	2600°K
Pressure	.. .. .	0.8 Atm	2.6 Atm.

No data for the spectra of the exhaust flames from rockets is available, but a reproduction of the spectrum of the exhaust flame from a jet engine is given in Fig. 5.

The spectrum consists of the well known bands of  $C_2$ , CH, and OH. Besides these, there are HCO bands which were first discovered by the author in 1934. A striking feature of the spectrum is the enhanced intensity of the HCO bands. (3B, Fig. 5).

One important aspect of the study of the spectra of the exhaust flames is to understand the mechanism of production of the different radicals, which throws interesting light on the processes of combustion. Such work is in progress in the Optics Division of the National Physical Laboratory.