Effect of Fe on the Martensitic Transition, Magnetic and Magnetocaloric Properties in Ni-Mn-In Melt-spun Ribbons

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ABSTRACT

The effect of Fe on the martensitic transitions, magnetic and inverse magnetocaloric effect in Ni$_{47}$Mn$_{40-x}$Fe$_{x}$In$_{13}$ ribbons (x = 1, 2, 3, and 5) has been investigated. All the ribbon compositions under study have shown the presence of austenite phase at room temperature. The variation of martensitic transition with the increase in Fe-content is non-monotonic. The thermal hysteresis of the martensitic transition increased with the increase in Fe-content. The martensitic transitions shifted to lower temperatures in the presence of high magnetic fields. A maximum magnetic entropy change ($\Delta S_m$) of 50 Jkg$^{-1}$K$^{-1}$ has been achieved in the Ni$_{47}$Mn$_{40}$Fe$_{1}$In$_{13}$ (x = 1) ribbon at 282 K for an applied field of 5 T.

Keywords: Thermo-magnetic behaviour, magnetocaloric effect, martensitic transition

1. INTRODUCTION

The Ni-Mn-Y (Y = Ga, In, Sn, and Sb) based ferromagnetic Heusler alloys display first order martensitic transitions from cubic to tetragonal or orthorhombic structure upon cooling. Associated with the martensitic transition observed in these alloys several important functional properties such as shape memory effect, inverse magnetocaloric effect and magnetoresistance have been reported$^{1-3}$. The martensitic and magnetic transitions in these alloys are composition dependent as they are highly sensitive to the valence electron concentration per atom (e/a) and Mn-Mn interatomic distances$^{4,5}$. In the case of Indium (In) containing alloys, a large inverse magnetocaloric effect (IMCE) in polycrystalline state and significant magnetic field induced strain (MFIS) in single crystals have been reported$^{6,7}$. Recently rapid quenching or melt spinning technique was attempted in various systems such as Ni-Fe-Ga, Ni-Mn-Y (Y = In, Al, Sn, and Ga) alloys$^{8-11}$. The melt spinning technique has the advantage of producing strong textured polycrystalline ribbons. In the case of magnetic refrigeration applications the use of refrigerant materials in the form of thin films or ribbons can optimize the heat transfer between the working body and the heat-exchange fluid. Additional advantage of melt spun ribbons is that the thin ribbons reduces the eddy current losses and can be utilised at high frequencies. In the recent past, off-stoichiometric Ni-Mn-In ribbons has gained lot of interest due to their intriguing martensitic phase transitions, microstructure, associated magnetic, magnetocaloric properties$^{12-17}$. In the present study, off-stoichiometric Ni$_{47}$Mn$_{40}$In$_{13}$, ribbon composition which reported a large magnetic entropy change of 32 Jkg$^{-1}$K$^{-1}$ at RT was chosen and Mn was partially substituted by Fe$^{18}$. The objective of this work is to understand the effect of Fe on the martensitic transitions, magnetic and magnetocaloric properties in Ni-Mn-In ribbons. Hence, Ni$_{47}$Mn$_{40-x}$Fe$_{x}$In$_{13}$ (x = 1, 2, 3, and 5) ribbons were prepared through melt spinning route and studied.

2. EXPERIMENTAL WORK

Four Ni-Mn(Fe)-In alloys were prepared by arc-melting pure metals in an argon atmosphere. Subsequently they were induction melted in a quartz tube and ejected with a 1 bar pressure difference onto a copper wheel rotating at a surface velocity of 17 m/s. The process was carried out in argon environment. As-quenched ribbons were 1.5 mm - 2 mm wide, 5 mm -18 mm long and 40 µm - 45 µm thick. As-spun ribbons were annealed at 850 °C for 30 min in vacuum and quenched to room temperature using Argon gas. The crystal structure of the constituent phases were examined using an X-ray diffractometer (PHILIPS3121, Cu-K$_\alpha$ λ = 1.54058 Å). The compositions of the ribbons were determined by energy dispersive x-ray spectroscopy (EDX) and were found to be close to the nominal composition. The phase transitions between the martensite and austenite phases were determined by differential scanning calorimetry (DSC) with a cooling/heating rate of 20 K/min. The magnetisation measurements were performed by means of vibrating sample magnetometer in physical property measurement system (PPMS, Quantum design P525) in the
temperature range of 4 K - 320 K. The sample was initially cooled in the absence of an external magnetic field and the data was collected on warming by applying an external magnetic field of 0.05 T termed as zero field cooled (ZFC) M(T) curve. Subsequently, the data was recorded upon cooling termed as field cooled (FC) M(T) curve without removing the applied field. The isothermal magnetisation curves were recorded for both increasing and decreasing fields for a maximum field of 5 T in the vicinity of structural transitions.

3. RESULTS AND DISCUSSION

3.1 Structural Analysis

Figure 1 (a)-(d) shows the x-ray diffraction patterns of Ni<sub>x</sub>Mn<sub>40-x</sub>Fe<sub>x</sub>In<sub>13</sub> ribbons recorded at room temperature (RT). The X-ray diffraction patterns of the x = 2 and 3 ribbons have been re-recorded using a step scan of 0.01° for a time interval of 4 s in the 2θ range of 30°-52.5° and 62.5°-80° (Fig. 1(e) and 1(f)). The x = 1, 2, and 5 ribbons revealed the presence of cubic L2<sub>1</sub> structure (austenite phase) at RT and no secondary phases were detected. This suggests that the martensitic transitions of these ribbons are below RT. In the case of x = 3 ribbon, step scan revealed the presence of a small volume fraction of the martensite along with major austenite phase at RT and no secondary phases were detected. This suggests that the martensitic transition of the x = 3 ribbon is around RT. The lattice parameters of the austenite phase in x = 1, 2, 3, and 5 ribbons are 0.5992 nm, 0.5987 nm, 0.5982 nm, and 0.5979 nm respectively. The partial substitution of Fe for Mn makes the lattice shrink because of the smaller atomic radius of Fe than that of Mn. Y. Feng et al. reported the presence of γ secondary phase in the Ni<sub>x</sub>Mn<sub>40-x</sub>Fe<sub>x</sub>In<sub>13</sub> alloys for y ≥ 5. This was found to be due to the solid solubility limit of Fe

in Ni-Mn-In alloys about 4 at. per cent. In the current ribbon series, no secondary phases were observed after doping Fe up to 5 at. per cent. This suggests that melt spinning technique has enhanced the solubility limit.

3.2 Phase Transformations

The cyclic DSC thermograms of Ni<sub>x</sub>Mn<sub>40-x</sub>Fe<sub>x</sub>In<sub>13</sub> melt spun ribbons (Fig. 2) display an exothermic peak for the phase transformation from austenite to martensite (A→M) upon cooling and reverses (M→A) as an endothermic peak upon heating. The values of Ms, Mf, As and Af are given in Table 1. Evidently, these ribbons display both first order structural transformation associated with martensitic transition and a second order magnetic transition (slope change in base line) corresponding to the Curie temperature of the austenite phase (T<sub>CT</sub>). All the ribbons except x = 2 show one peak which corresponds to the structural transition (M→A) followed by a magnetic transition of the austenite phase (T<sub>CT</sub>) upon heating. The x = 2 ribbon shows two peaks with significant thermal hysteresis upon heating and cooling suggesting the presence of two structural transitions. In the cooling run, the first exothermic peak corresponds to the premartensitic transition (PMT) ~ 294 K followed by a martensitic transition ~ 255 K. This PMT is believed to be the precursor for the martensitic transition. The magneto-elastic coupling was found to play an important role for the occurrence of this precursor effect. Very recently, it was reported that PMT was observed in Sn-doped Ni<sub>50</sub>Mn<sub>34</sub>Ga at the Ga site up to 5 at. per cent<sup>20</sup>. This suggests that the presence of Fe influences the magneto-elastic coupling in Ni-Mn-In ribbons. Except in the x = 2 ribbon sample, PMT was not observed in other compositions which could be due to other factors suppressing PMT which needs further exploration.

Figure 1. X-ray diffraction patterns for Ni<sub>x</sub>Mn<sub>40-x</sub>Fe<sub>x</sub>In<sub>13</sub> (a) x = 1 (b) x = 2 (c) x = 3 (d) x = 5 (e) and (f) slow scans of x = 2 and 3 ribbons in selected ranges of 2θ respectively at RT.
Figure 2. DSC curves for Ni$_{40}$Mn$_{40-x}$Fe$_x$In$_{13}$ ($x = 1, 2, 3, 5$) ribbons upon heating and cooling at a rate of 20 K/min.

Fe$_x$In$_{13}$ ribbons carried out as a function of temperature in a field of 0.05 T (500 Oe) are as shown in Fig. 3. All the ribbons ($x = 1-5$) show Curie transitions of the martensite phases ($T_M^C$) at low temperatures and splitting between the FC and ZFC curves below $T_M^C$. The $T_M^C$ of the $x = 1, 2$, and 3 ribbons were observed to be overlapping with its martensitic phase transition. The $T_M^C$ of the base composition Ni$_{47}$Mn$_{40}$In$_{13}$ ribbon was ~120 K and it has increased to 165 K for substituting Fe up to 3 at.% for Mn. With further increase in Fe (5 at.%), $T_M^C$ has increased to ~200 K. This suggests that the presence of Fe enhances the ferromagnetic exchange interactions in the martensite phase in this series. At low temperatures (T<100 K), another transition $T_{EB}^M$ called exchange bias blocking temperature, is observed in the ZFC curves. The observation of $T_{EB}^M$ and splitting of FC and ZFC curves suggest that there exist antiferromagnetic (AFM) and ferromagnetic (FM) interactions in these ribbons at low temperatures. The exchange bias blocking temperatures ($T_{EB}^M$) of the $x = 1, 2, 3, 5$ and 5 ribbons are 55 K, 25 K, 51 K, and 7 K respectively. The variation in the $T_{EB}^M$ values with the increase in Fe-content was non-monotonous, but the general trend (except for the $x = 3$ ribbon) appears to be decreasing with increase in Fe-content. This may be attributed to the decrease in the FM-AFM coupling strength with the increase in Fe-content as the ferromagnetic exchange in the martensite phase was observed to be enhanced with the increase in Fe-content.

From the dM/dT curves upon heating and cooling shown in the insets, thermal hysteresis of the martensitic transitions were found to increase with increase in Fe-content which is in good agreement with the DSC results. These observations are similar to that reported in NiMnFeGa alloys which convey that the presence of Fe atoms in the NiMnGa systems increases the energy required for phase boundary motion which might result in a decrease of the thermal elasticity in martensitic understanding. It was also found that the thermal hysteresis increases with the increase in Fe-content. For the $x = 1$ ribbon, $T_M$ was around 292 K and for $x = 2$ ribbon, the $T_M$ lowered to 270 K. In the case of $x = 3$ ribbon, the $T_M$ increased to 303 K and again drastically lowered to 218 K for the $x = 5$ ribbon.

The X-ray analysis of the $x = 3$ ribbon revealed the presence of low volume fraction of the martensite phase along with major austenite phase at RT. The DSC results support the XRD studies. It can be seen that the variation of $T_M$ with the increase in Fe concentration is non-monotonic. $T_M$ is linearly related to the valence electron concentration ($e/a$) of the ribbon. Therefore, it was observed that the variation of $T_M$ with composition is expected due to the monotonic change in the number of valence electrons which has been well reported in various analogous systems. However, in the present case the variation of $T_M$ with $e/a$ was observed to be non-linear. This suggests that additional parameters other than electron concentration could affect phase stability. The non-monotonous nature of $T_M$ as a function of $e/a$ in the present study suggests that the effect of alloying is not just a change in the Fermi level, but the addition of Fe could also modify to some extent the orbital hybridisation and bonding. The changes in hybridisation and related volume effects were reported for Ni$_{47}$Mn$_{40}$Ga with substitutions. Hence, a simple choice of $e/a$ can only be a guideline for examining systematic changes within a single alloy system and is not universal for related quaternary systems. It was reported that the dopant would create negative/positive pressure due to expansion/contraction of lattice volume. Such pressure effect can dominate over the effect related to $e/a$. It has been observed that hydrostatic pressure can enhance the MT in many Ni-Mn based alloys. In the current study, similar effects are possible with the Fe substitution in Ni-Mn-In ribbons. Hence, the influence of multiple factors like composition (or $e/a$), pressure (or volume) and atomic disorder led to the observed complex behaviour of $T_M$ with Fe-content.

The ZFC and FC magnetisation measurements of Ni$_{47}$Mn$_{40}$, Mn$_{40}$In$_{13}$.
transformation. In the case of a Curie transition, the $dM/dT$ curves should overlap upon heating and cooling indicating no thermal hysteresis as it is a second order transition. But, in this alloy series, the cyclic $dM/dT$ curves of $x = 2, 3, and 5$ ribbons have also shown small thermal hysteresis (referring to the negative $dM/dT$ peaks) in the vicinity of their respective Curie transitions unlike that observed in the case of second order transitions. This can be attributed to the presence of martensitic transition in the vicinity of Curie transition ($T_{C}^{M}$) that is giving rise to the observed thermal hysteresis.

3.3 Exchange Bias

A shift in field cooled M-H curve along field axis is the characteristic behaviour of exchange bias phenomena. It was verified by measuring the ZFC and FC hysteresis loops of $x = 1$ ribbon at different temperatures (a) 5 K, (b) 30 K, (c) 50 K, (d) 70 K (Fig. 4) in the field range of -2 T to 2 T. However, for better clarity the loops were shown from -0.2 T to 0.2 T in Fig. 4. It was found that, the FC loop shifts to negative field from origin up to a temperature of 50 K and at 70 K the shift was absent. On the other hand, the ZFC loop did not show any shift for all the set temperatures in the range of 5 K – 70 K. It implies that the EB phenomenon present for temperatures below $T_{EB}^{c}$ (~ 55 K) for the $x = 1$ ribbon. The coercivity ($H_{c}$) and EB field ($H_{E}$) values were determined from the FC hysteresis loops using $H_{E} = -(H_{1}+H_{2})/2$ and coercivity $H_{c} = |H_{1}-H_{2}|/2$ where the $H_{1}$ and $H_{2}$ are the negative field and positive field at which the magnetisation equals to zero. On increasing the temperature the $H_{E}$ (Fig. 5) is found to decrease linearly and become zero as the temperature approaches $T_{EB}^{c}$ (~ 70 K), whereas the $H_{c}$ increases up to a maximum value of 260 Oe at 70 K and after that it decreases with further increase in temperature. The disappearance of $H_{E}$ above $T_{EB}^{c}$ confirms that the EB phenomenon exists only below $T_{EB}^{c}$. The $H_{E}$ and $H_{c}$ values are 175 Oe and 55 Oe at 10 K, respectively. The $H_{c}$ curve presents a maximum value of 260 Oe at 70 K. Such dependencies of $H_{E}$ and $H_{c}$ are typical for EB systems. The weakening of AFM-FM coupling due to the decrease of AFM region size on increasing temperature could be the reason for the observed trend in $H_{E}$ and $H_{c}$. The evidence of EB results in Ni$_{47}$Mn$_{40-x}$Fe$_{x}$In$_{13}$ ribbons signifies the existence of AFM-FM interactions. However, the other methods such as neutron diffractions studies are required to probe these kinds of magnetic structures in further detail.

3.4 Effect of Magnetic Field on the Martensitic Transitions

To understand the effect of magnetic field on structural and magnetic transitions, the thermomagnetic curves of Ni$_{47}$Mn$_{40-x}$Fe$_{x}$In$_{13}$ ribbons were measured in a high magnetic field of 5 T.
upon heating and cooling in the temperature range of 5 K - 330 K (Fig. 6). The features in the high field curves are also similar to those shown in low field data (Fig. 3), but exhibit broad transformations. These ribbons show a shift in martensite start temperatures \( M_s \) towards lower temperatures as the applied magnetic field increases (Table 2).

For the \( x = 1, 2, 3 \) and 5 ribbons, the \( M_s \) values shifted by 15 K, 25 K, 13 K, and 66 K respectively, for a field change of 4.95 T. The \( x = 5 \) ribbon showed a maximum rate of change of \( M_s \) with respect to field (11.2 K/ T) in this alloy series. The change of \( M_s \) with field signifies a field induced transition from the martensite to the austenite phase. The lowering of the martensitic transition in these ribbons implies that the magnetic field favors the formation of the austenitic phase. Conversely, in Ni-Mn-Ga alloys, the \( M_s \) was reported to shift to higher temperatures with increasing magnetic field\textsuperscript{28}. This kind of behaviour is due to the fact that the saturation magnetisation of the martensitic phase is larger than that of the austenite in Ni–Mn–Ga in contrast to that of Ni–Mn–X (X = Sn, In, Sb), wherein the martensitic phase has a lower magnetisation than

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**Figure 4.** The ZFC and FC hysteresis loops of Ni\(_{47}\)Mn\(_{40-x}\)Fe\(_x\)In\(_{13}\) ribbon \( (x = 1) \) (a) 10 K, (b) 30 K, (c) 50 K, (d) 70 K. The sample is cooled in a field of 5 T for FC loops. The loops are displayed from -0.2 to 0.2 T for clear visualisation of the loop shift.

**Figure 5.** Exchange bias field \( (H_e) \) and coercivity \( (H_c) \) as a function of temperature in Ni\(_{47}\)Mn\(_{40-x}\)Fe\(_x\)In\(_{13}\) ribbon \( (x = 1) \).

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**Table 2.** Martensite start temperatures \((M_s)\) at 50 mT and 5 T fields and the shift in the martensite start \( M_s (\Delta T) \) both experimental and calculated from Clausius-Clapeyron relation

<table>
<thead>
<tr>
<th>( x )</th>
<th>( M_s ) (50 mT) (K)</th>
<th>( M_s ) (5 T) (K)</th>
<th>( \Delta T ) (K) [Expt.]</th>
<th>( \Delta T ) (K) [C-C equ.]</th>
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<td>5</td>
<td>208</td>
<td>142</td>
<td>66</td>
<td>20</td>
</tr>
</tbody>
</table>

For the \( x = 1, 2, 3 \) and 5 ribbons, the \( M_s \) values shifted by 15 K, 25 K, 13 K, and 66 K respectively, for a field change of 4.95 T. The \( x = 5 \) ribbon showed a maximum rate of change of \( M_s \) with respect to field (11.2 K/ T) in this alloy series. The change of \( M_s \) with field signifies a field induced transition from the martensite to the austenite phase. The lowering of the martensitic transition in these ribbons implies that the magnetic field favors the formation of the austenitic phase. Conversely, in Ni-Mn-Ga alloys, the \( M_s \) was reported to shift to higher temperatures with increasing magnetic field\textsuperscript{28}. This kind of behaviour is due to the fact that the saturation magnetisation of the martensitic phase is larger than that of the austenite in Ni–Mn–Ga in contrast to that of Ni–Mn–X (X = Sn, In, Sb), wherein the martensitic phase has a lower magnetisation than
the austenitic phase. The \( M_s \) shift (\( \Delta T \)) induced by magnetic field change \( \Delta H \) can also be calculated approximately by Clausius-Claypeyron relation [C-C relation]

\[
\Delta T = \left( \frac{\Delta M}{\Delta S} \right) \Delta H
\]

where \( \Delta M \) and \( \Delta S \) are the differences in magnetisation and thermal entropy between the austenite and martensite phases. The \( \Delta S \) values of the ribbons are calculated from the enthalpy data obtained from DSC. In the case of \( x = 2 \) ribbon where two transitions are observed, the sum of the two thermal entropies at those transitions have been taken into account for calculation. The experimental and the theoretical values (C-C equation) of the shift in the \( M_s \) are comparable in the case of \( x = 1, 2 \) and \( 3 \) ribbons. In the case of \( x = 5 \) ribbon, the \( M_s \) values do not match because some of the volume fraction of the austenite phase is retained even at low temperatures (up to 4 K) in the presence of high field. Hence, the C-C equation is valid only when the phase transformation from austenite to martensite phase is complete.

A kinetic arrest phenomenon has been reported in Ni-Mn-In alloys\(^{29,30} \). In this phenomenon, the MT is interrupted at certain temperatures during magnetic field cooling and does not proceed with further cooling. The parent (austenite) phase remains down to low temperatures. From the FC and FH \( M(T) \) curves of the \( x = 5 \) ribbon, it can be said that this ribbon displays kinetic arrest phenomenon because of which the transformation to martensite upon cooling was not complete. Although such a peculiar behaviour of entropy change would be associated with the magnetic contribution to the Gibbs energy of the parent phase, the origin is still under discussion.

### 3.5 Magnetic Isotherms and Magnetocaloric Effect

The isothermal magnetic entropy change (\( \Delta S_M \)) as a function of temperature was studied for the Ni\(_{47}\)Mn\(_{40-x}\)Fe\(_x\)In\(_{13}\) ribbons by measuring the isothermal magnetisation curves at different temperatures in the vicinity of their respective martensitic transformations. The isothermal magnetisation curves were measured in both increasing (0-5 T) and decreasing (5-0 T) fields. The isotherms of the ribbons were measured at a temperature interval of 2 K around \( T_M \). The typical cyclic M-H loops of the \( x = 1 \) and 3 ribbons (Fig. 7) exhibit magnetic hysteresis around their respective \( T_M \) due to the magnetic field induced reverse martensitic transformations. The magnetisation values were found to increase with increasing temperature in the ranges of 276 K - 294 K and 292 K - 304 K for the \( x = 1 \) and 3 ribbons respectively, in the austenite region and upon further increase in temperature magnetisation decreases in both the cases. This is due to the transformation of martensite phase into austenite phase. Similar behaviour was observed in the
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The critical field \( H_{cr} \) defining the transition point is determined as the crossing point of the linear portion of the curves. This behaviour is attributed to the onset of magnetic field induced structural transition (FIST) from martensite to austenite phase. The \( x = 1 \) and \( x = 2 \) ribbons have shown the presence of strong FIST, whereas, the \( x = 3 \) ribbon showed feeble FIST and in the case of \( x = 5 \), the FIST was observed to be moderate. In the case of \( x = 1, 2, \) and \( 5 \) ribbons, the shift in \( H_{cr} \) to lower fields (shown by solid line with arrow) with increasing temperature could be observed. This is because the martensite to austenite phase transition is driven by both field and temperature, the critical field required to drive the structural transition at higher temperature in the magneto-structural transition regime is smaller. The absence of a strong FIST in the \( x = 3 \) ribbon suggests that the critical field required to induce a strong FIST is more than 5 T. Temperature alone is enough for driving the structural transition and in addition to that if magnetic field is applied, it becomes much easier to induce a strong FIST at lower critical fields. Apart from \( H_{cr} \), there is another slope change observed in the M-H curves upon increasing fields (shown by dotted line) around 4.5 T. This can be explained as follows: the maximum applied field (here 5 T) is not enough to reach saturation point and upon decreasing the field, the M-H curve will reverse with higher magnetisation value. The M-H curve upon increasing field is only getting connected to the M-H curve upon decreasing field resulting in the slope change.

The temperature dependence of \( \Delta S_m \) calculated in the decreasing field for \( \text{Ni}_{47}\text{Mn}_{40-x}\text{Fe}_x\text{In}_{13} \) ribbons are shown in Fig. 6. A maximum positive \( \Delta S_m \) value (IMCE) of 50 J/kg K has been achieved both in the \( x = 1 \) and \( 2 \) ribbons, however at different temperatures ~ 282 K and 255 K respectively for a \( \Delta H \) of 5 T. These values are more than that reported in \( \text{Ni}_{41.9}\text{Fe}_{y}\text{Mn}_{37}\text{In}_{13} \) \( (y = 1, 3) \) alloys\(^\text{11} \). On the other hand, maximum positive \( \Delta S_m \) values of ~35 J/kg K at 298 K and ~25 J/kg K at 208 K have been obtained in the \( x = 3 \) and \( 5 \) ribbons respectively. The refrigerant capacities (RC) of the ribbons have also been summarised in the Table 3.

The obtained large \( \Delta S_m \) values around \( T_m \) in the \( x = 1 \) and \( 2 \) ribbons is due to large change in the magnetisation values of

\[ \Delta S_m = \text{Temperature dependence of } \Delta S_m \text{ in } \text{Ni}_{47}\text{Mn}_{40-x}\text{Fe}_x\text{In}_{13} \text{ ribbons measured on decreasing field (5-0 T).} \]

(a) \( x = 1 \), (b) \( x = 2 \), (c) \( x = 3 \) and (d) \( x = 5 \).
the martensite and austenite phases in a narrow temperature interval arising due to a strong field induced structural transition compared to those observed in x = 3 and 5 ribbons. The $\Delta S_M$ values decrease with increase in Fe-content suggesting that Fe may aid in suppressing the field induced structural transition. The refrigerant capacity (RC) values for the Ni$_{47}$Mn$_{40-x}$FexIn$_{13}$ ribbons calculated from the integral of $\Delta S_M$ over the full width at half maxima are given in Table 3. The RC values are estimated to be 227 J/kg, 204 J/kg, 140 J/kg, and 202 J/kg on decreasing fields for x = 1, 2, 3, and 5 ribbons, respectively. The estimated average hysteresis losses for the ribbons at temperatures where $\Delta S_M$ values attain their respective maxima are also given in the Table 3. Therefore, the net RC values of the ribbons obtained by subtracting the average hysteresis losses from their RC values were 117 J/kg, 116 J/kg, 88 J/kg, and 202 J/kg on (0-5 T) for x = 1, 2, 3, and 5, respectively.

### Table 3. Values of magnetic entropy change ($\Delta S_M$), temperature at which IMCE attains maxima ($T_{IMCE}^{\text{max}}$) and refrigeration capacity (RC) in Ni$_{47}$Mn$_{40-x}$FexIn$_{13}$ ribbons obtained on decreasing field in the vicinity of martensitic transitions

<table>
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<th>$\Delta S_M$ (J/kg K) (5-0 T)</th>
<th>$T_{IMCE}^{\text{max}}$ (K)</th>
<th>RC (J/kg) (5-0 T)</th>
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4. **Summary**

Structure, thermo-magnetic behaviour and magnetocaloric effect have been studied in Ni$_{47}$Mn$_{40-x}$FexIn$_{13}$ ribbons. All the compositions (x = 1 to 5) have shown the presence of austenite phase at RT. It was found that the variation of martensitic transition ($T_m$) with the increase in Fe (x = 1 to 5) concentration was non-monotonic. This suggests that the exclusive dependence of martensitic transitions on e/a values is not universal and in some compositions (like in the present case) other factors are predominant. The presence of Fe did not have much influence on the $T_m$. The thermal hysteresis of the martensitic transition increased with the increase in Fe-content. The ribbons showed a shift in martensite start temperatures towards lower temperatures in the presence of high magnetic fields.

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Large magnetic entropy change in Ni$_{50-x}$Mn$_{34}$In$_x$ Heusler alloys. *Appl. Phys. Lett.*, 2007, **90**, 262504. doi: 10.1063/1.2752720


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**CONTRIBUTORS**

Dr. D M Raj Kumar received M Tech (Solid State Technology) from IIT Kharagpur and PhD (Physics) from IIT Bombay. He is working as Scientist in the advanced magnetics group, DMRL, Hyderabad. His area of research interests include: Rare-earth permanent magnets and magnetocaloric materials. In the current study, his contribution: Preparation of melt spun ribbons, structural characterisation and complete analysis

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